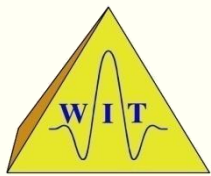


3D i-CRS stacking operator

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University of Hamburg, St. Petersburg University

Wave Inversion Technology



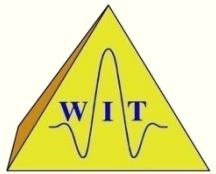
2D i-CRS stacking operator (Schwarz et al., 2014):

- Double-square-root expression
- Robust regarding heterogeneity
- Leads to good results for high reflector curvatures
- Provides higher accuracy than CRS and conventional MF

Requirements for 3D i-CRS stacking operator:

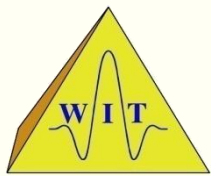
- Computationally efficient solution
- Expressed in CRS parameters

Outline



- Introduction
- Theory:
 - $2D + ? = 3D$
 - Why, why, how?
 - Make it easier!
- Quantitative studies
- Synthetic data example
- Conclusions and outlook

Introduction: 2D i-CRS



- CRS formula in 2D case reads:

$$t^2(\Delta x_m, h) = \left[t_0 + w \Delta x_m \right]^2 + N \Delta x_m^2 + M h^2$$

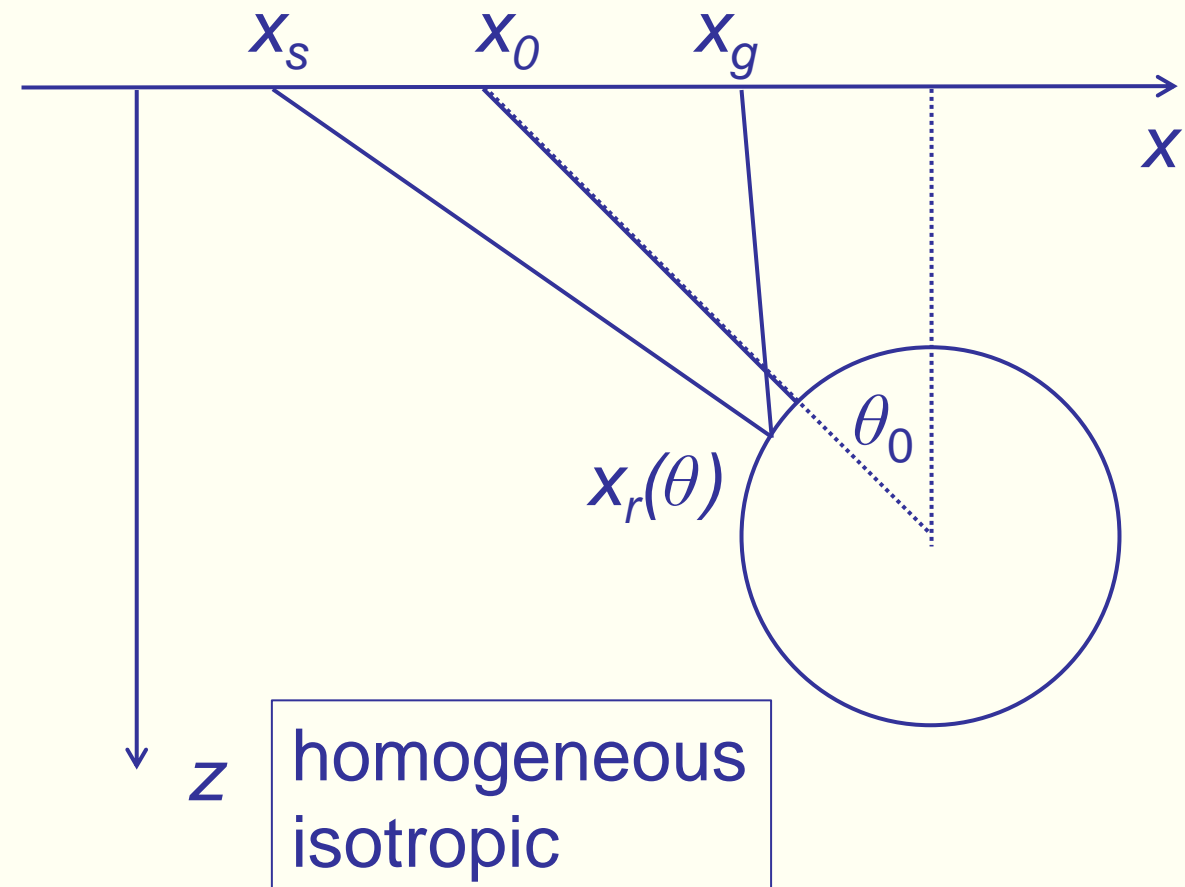
- DSR expression for travelttime:

$$t = \sqrt{|\mathbf{x}_s - \mathbf{x}_r|^2} / V + \sqrt{|\mathbf{x}_g - \mathbf{x}_r|^2} / V$$

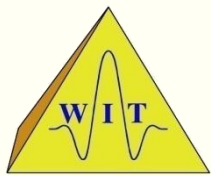
- Reflection point:

$$\frac{\partial t(\theta)}{\partial \theta} = 0,$$

$$\tan \theta = \tan \theta_0 + f(t_s, t_g)$$



Theory: 2D + ? = 3D



- CRS formula in 3D case reads:

$$t^2(\Delta \mathbf{x}_m, \mathbf{h}) = \left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_m \right]^2 + \Delta \mathbf{x}_m^T \mathbf{N} \Delta \mathbf{x}_m + \mathbf{h}^T \mathbf{M} \mathbf{h}$$

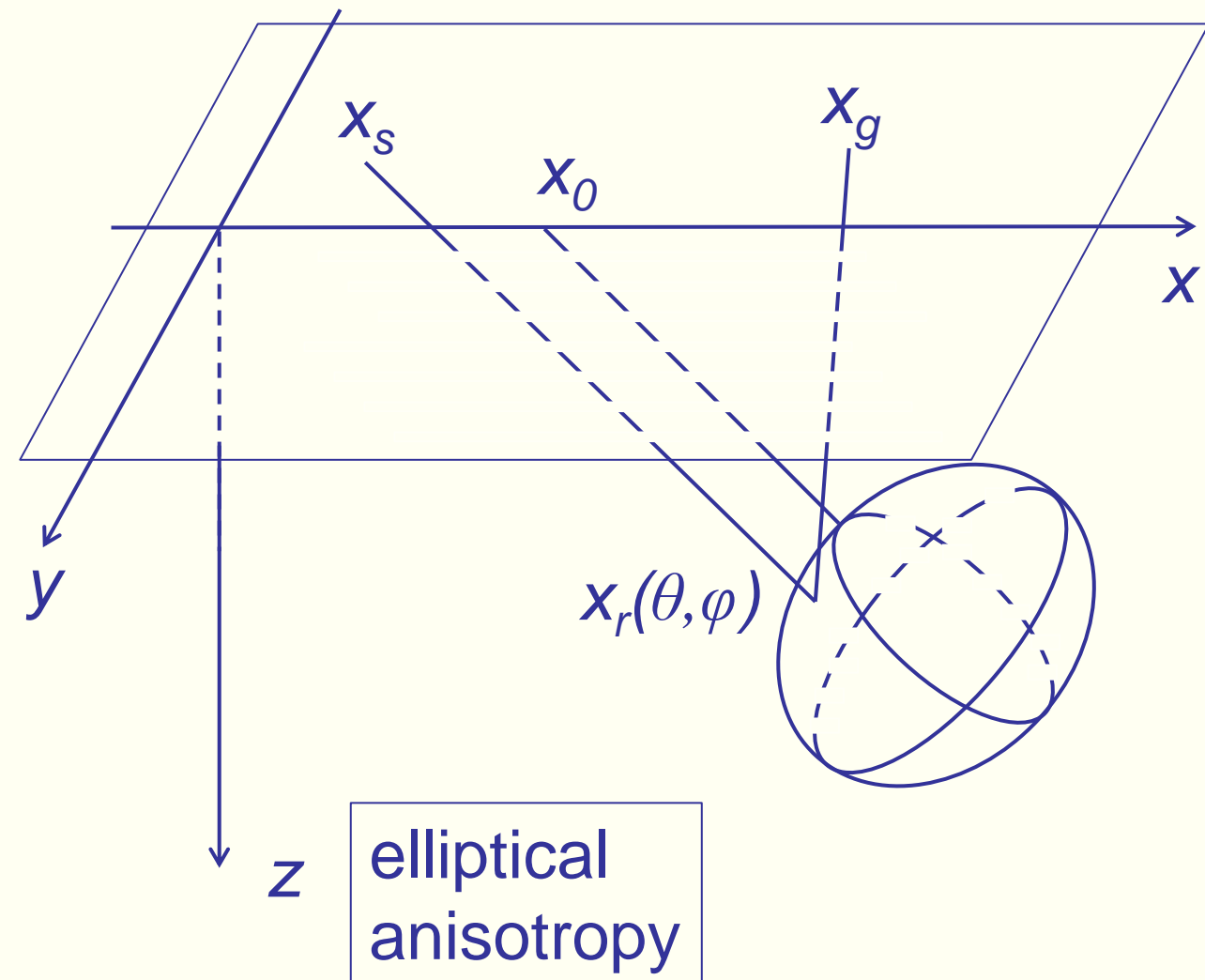
- DSR expression for traveltimes:

$$t = \sqrt{|\mathbf{x}_s - \mathbf{x}_r|^2} / V_s + \sqrt{|\mathbf{x}_g - \mathbf{x}_r|^2} / V_g$$

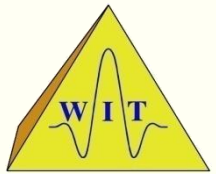
- Reflection point:

$$\frac{\partial t(\theta, \phi)}{\partial \theta} = 0,$$

$$\frac{\partial t(\theta, \phi)}{\partial \phi} = 0$$

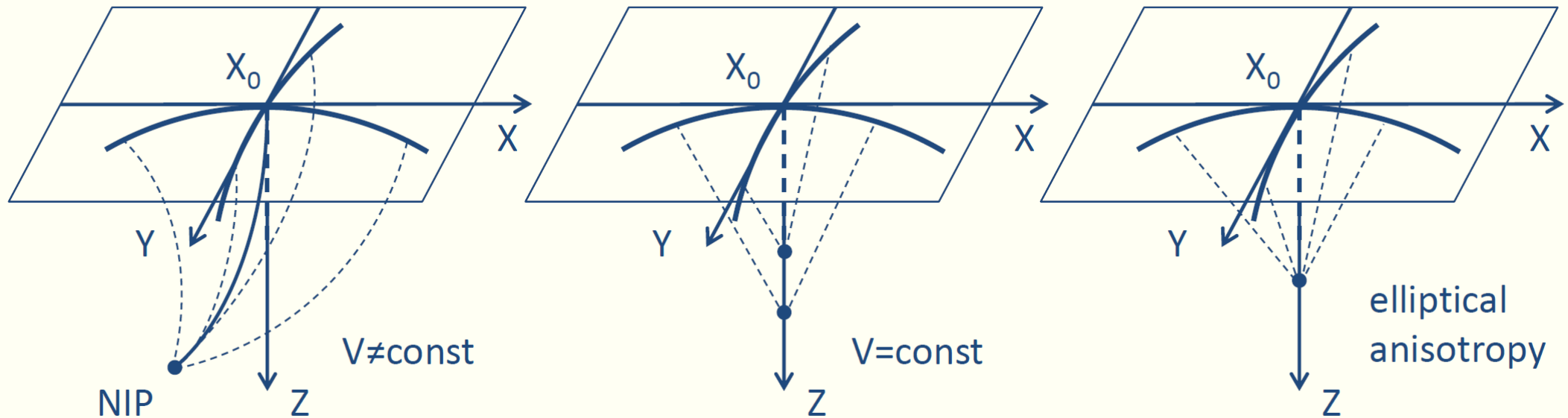
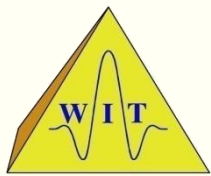


Why, why, how?



- Why is the auxiliary elliptical anisotropic medium required?
- Why is the ellipsoid considered as reflector?
- How does the presented algorithm work?

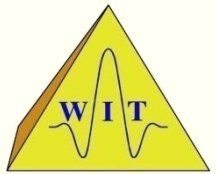
Why elliptical anisotropy?



- Elliptical anisotropic medium allows to focus wavefront with arbitrary curvature K_{NIP}
- Simple relationship between anisotropic parameters and curvatures of wavefront:

$$A_{11} = \frac{2}{t_0} \frac{\zeta_0}{K_{NIP}^{11}}, \quad A_{22} = \frac{2}{t_0} \frac{\zeta_0}{K_{NIP}^{22}}, \quad A_{33} = \zeta_0^2$$

Why ellipsoid?



Reflector may be locally approximated in vicinity of NIP point:

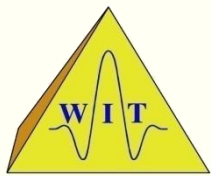
$$z = z_0 - \frac{1}{2} \begin{pmatrix} x \\ y \end{pmatrix}^T \mathbf{K} \begin{pmatrix} x \\ y \end{pmatrix}$$

Let us consider now ellipsoid:

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} + \frac{2xy}{c^2} + \frac{z^2}{d^2} = 1, \quad \mathbf{K} = \begin{pmatrix} d/a^2 & d/c^2 \\ d/c^2 & d/b^2 \end{pmatrix}$$

Note! There are many ellipsoids with curvature \mathbf{K} . There is one free parameter: d

How? 3D i-CRS stacking operator



Traveltime of reflected wave is:

$$t(\theta, \phi) = \frac{\sqrt{X_s^2 + Y_s^2 + Z_s^2}}{\zeta(\theta_s, \phi_s)} + \frac{\sqrt{X_g^2 + Y_g^2 + Z_g^2}}{\zeta(\theta_g, \phi_g)}$$

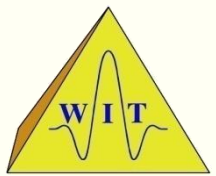
where angles:

$$\tan \phi_i = \frac{Y_i}{X_i}, \quad \tan \theta_i = \frac{\sqrt{X_i^2 + Y_i^2}}{Z_i}, \quad i = s, g$$

define group velocity:

$$\frac{1}{\zeta^2(\theta_i, \phi_i)} = \frac{\sin^2 \theta_i \cos^2 \phi_i}{A_{11}} + \frac{\sin^2 \theta_i \sin^2 \phi_i}{A_{22}} + \frac{\cos^2 \theta_i}{A_{33}}$$

How? 3D i-CRS stacking operator



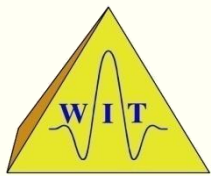
Travelttime of reflected wave is:

$$t(\theta, \phi) = \frac{\sqrt{\frac{A_{33}}{A_{11}} X_s^2 + \frac{A_{33}}{A_{22}} Y_s^2 + Z_s^2}}{\zeta_0} + \frac{\sqrt{\frac{A_{33}}{A_{11}} X_g^2 + \frac{A_{33}}{A_{22}} Y_g^2 + Z_g^2}}{\zeta_0}$$

To find the reflection point we should minimize the travelttime and therefore to calculate the derivatives:

$$\frac{\partial t(\theta, \phi)}{\partial \theta} = 0, \quad \frac{\partial t(\theta, \phi)}{\partial \phi} = 0$$

How? Iterative scheme



Condition $\frac{\partial t}{\partial \phi} = 0$ leads to the implicit equation for ϕ :

$$\tan \phi = \frac{A_{11} b \left(\frac{Y_s}{R_s} + \frac{Y_g}{R_g} \right)}{A_{22} a \left(\frac{X_s}{R_s} + \frac{X_g}{R_g} \right)} \quad (1)$$

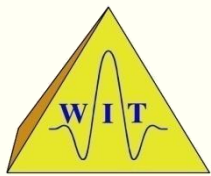
Condition $\frac{\partial t}{\partial \theta} = 0$ leads to the implicit equation for θ :

$$\tan \theta = \frac{\left(\frac{\Lambda_s}{R_s} + \frac{\Lambda_g}{R_g} \right)}{\left(\frac{Z_s}{R_s} + \frac{Z_g}{R_g} \right)} \quad (2)$$

where:

$$\Lambda_i = \frac{A_{33} a}{A_{11} d} X_i \cos \phi + \frac{A_{33} b}{A_{22} d} Y_i \sin \phi, \quad i = s, g$$

How? Iterative scheme



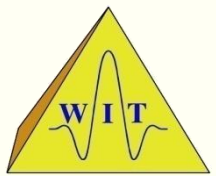
- Zero iteration: $\theta^{(0)} = 0$, $\phi^{(0)} = \arctan(y_m/x_m)$

- With initial values of angles we compute:

$$X_i^{(0)}, Y_i^{(0)}, Z_i^{(0)}, R_i^{(0)}, \quad i = s, g$$

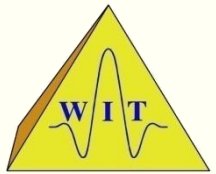
- Formula (1) updates $\phi^{(1)}$
- With $\phi^{(1)}$ we compute $\Lambda_s^{(0)}, \Lambda_g^{(0)}$
- Formula (2) updates $\theta^{(1)}$
- If required, we make next iteration
- Finally, after n iterations we obtain angles $\theta^{(n)}$ and $\phi^{(n)}$ and substitute them into DSR formula for traveltime

Make it easier!



- In case of diffracted waves ellipsoid collapses to the point

Make it easier!



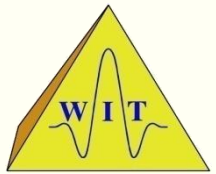
- The hyperbolic CRS operator:

$$t_{hyp}^2(\Delta \mathbf{x}_m, \mathbf{h}) = \left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_m \right]^2 + \Delta \mathbf{x}_m^T \mathbf{M} \Delta \mathbf{x}_m + \mathbf{h}^T \mathbf{M} \mathbf{h}$$

- i-CRS stacking operator, case of diffracted waves, exact formula:

$$t_{dif}(\Delta \mathbf{x}_m, \mathbf{h}) = \frac{1}{2} \sqrt{\left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_S \right]^2 + \Delta \mathbf{x}_S^T \mathbf{M} \Delta \mathbf{x}_S} + \frac{1}{2} \sqrt{\left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_G \right]^2 + \Delta \mathbf{x}_G^T \mathbf{M} \Delta \mathbf{x}_G}$$

Make it easier!

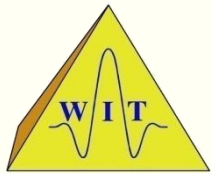


- Reflection point may be approximated by:

$$\Delta \mathbf{x}_r \approx [\mathbf{I} - \mathbf{K}_{NIP}^{-1} \mathbf{K}_N] \mathbf{R}^T \Delta \mathbf{x}_m$$

- \mathbf{R} is a matrix that accounts for the transformation from the ray-centred to the general Cartesian coordinate system
- \mathbf{K}_{NIP} and \mathbf{K}_N are curvatures of NIP and normal wave

Make it easier!



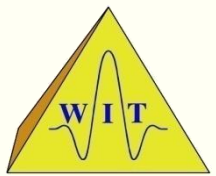
- The hyperbolic CRS operator:

$$t_{hyp}^2(\Delta \mathbf{x}_m, \mathbf{h}) = \left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_m \right]^2 + \Delta \mathbf{x}_m^T \mathbf{N} \Delta \mathbf{x}_m + \mathbf{h}^T \mathbf{M} \mathbf{h}$$

- i-CRS stacking operator, general case, approximation:

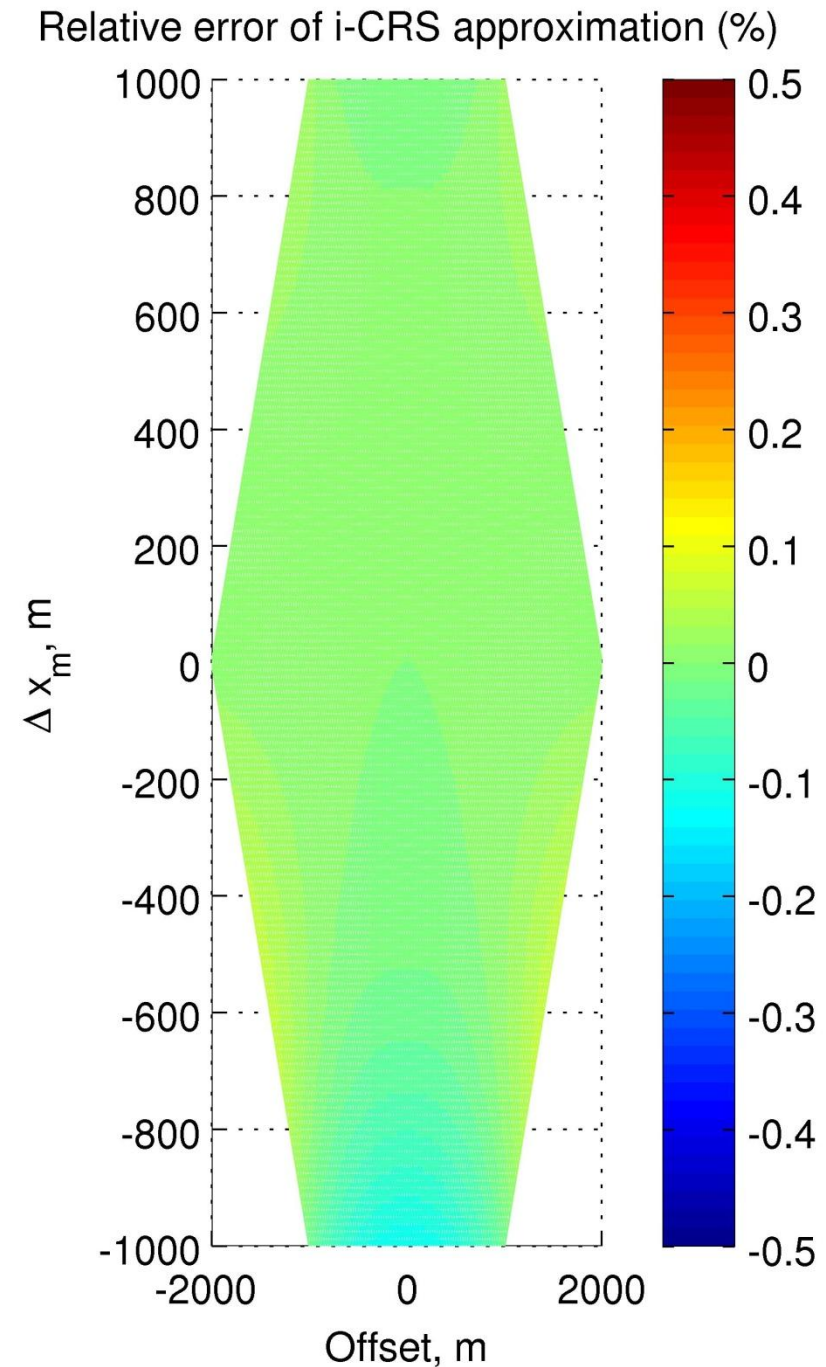
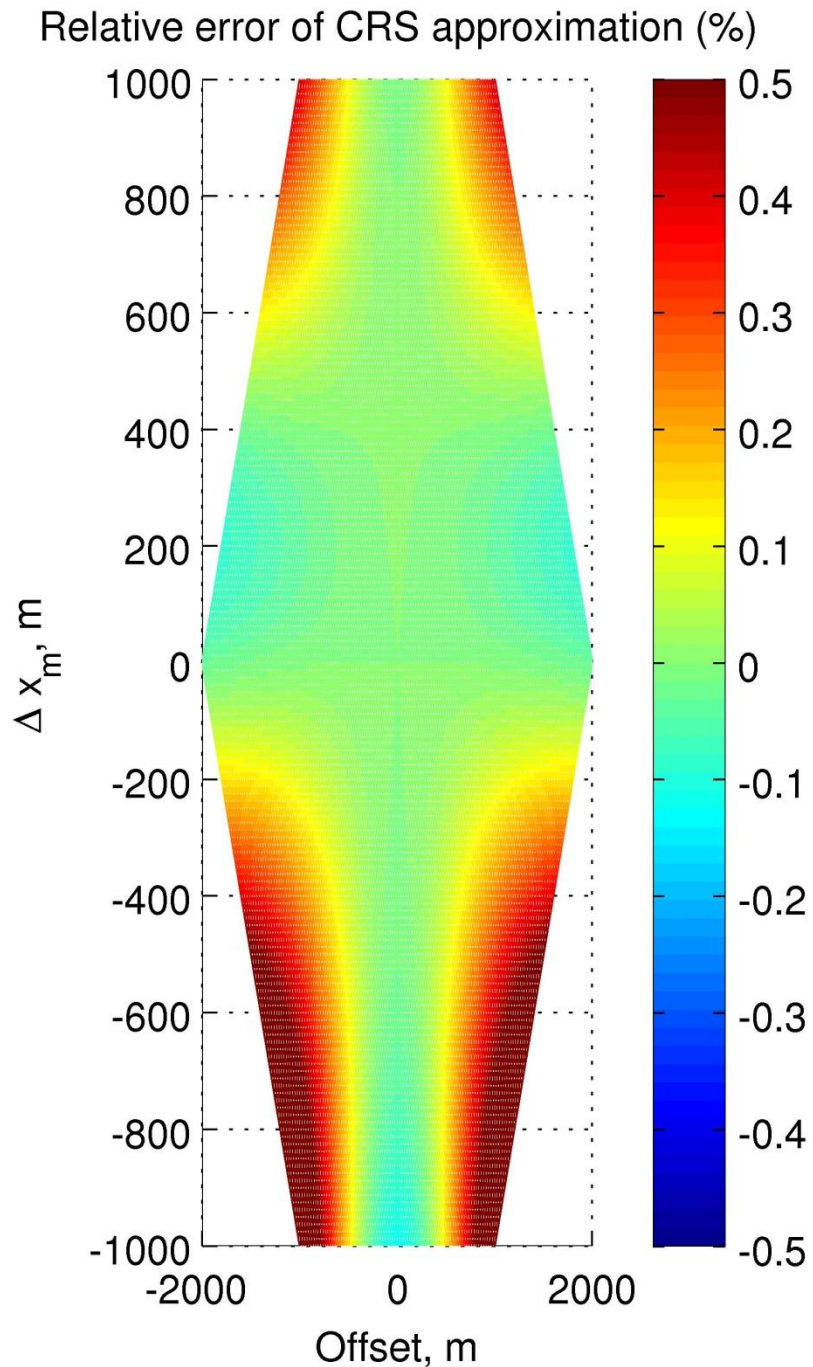
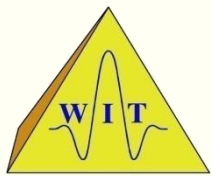
$$t_{ref}(\Delta \mathbf{x}_m, \mathbf{h}) =$$
$$\frac{1}{2} \sqrt{\left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_S \right]^2 + \Delta \mathbf{x}_m^T \mathbf{N} \Delta \mathbf{x}_m - 2 \Delta \mathbf{x}_m^T \mathbf{N} \mathbf{h} + \mathbf{h}^T \mathbf{M} \mathbf{h}}$$
$$+ \frac{1}{2} \sqrt{\left[t_0 + \mathbf{w}^T \cdot \Delta \mathbf{x}_G \right]^2 + \Delta \mathbf{x}_m^T \mathbf{N} \Delta \mathbf{x}_m + 2 \Delta \mathbf{x}_m^T \mathbf{N} \mathbf{h} + \mathbf{h}^T \mathbf{M} \mathbf{h}}$$

Numerical example: traveltimes

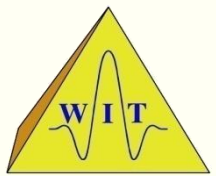


#	Reflector type	Radius	Medium	Depth
1	Sphere	1 km	Const	2 km
2	“Diffractor”	0.01 km	Const	2 km
3	“Plane”	10 km	Const	2 km
4	Ellipsoid	0.9, 2.8 km	Gradient	2 km

Numerical example: traveltimes

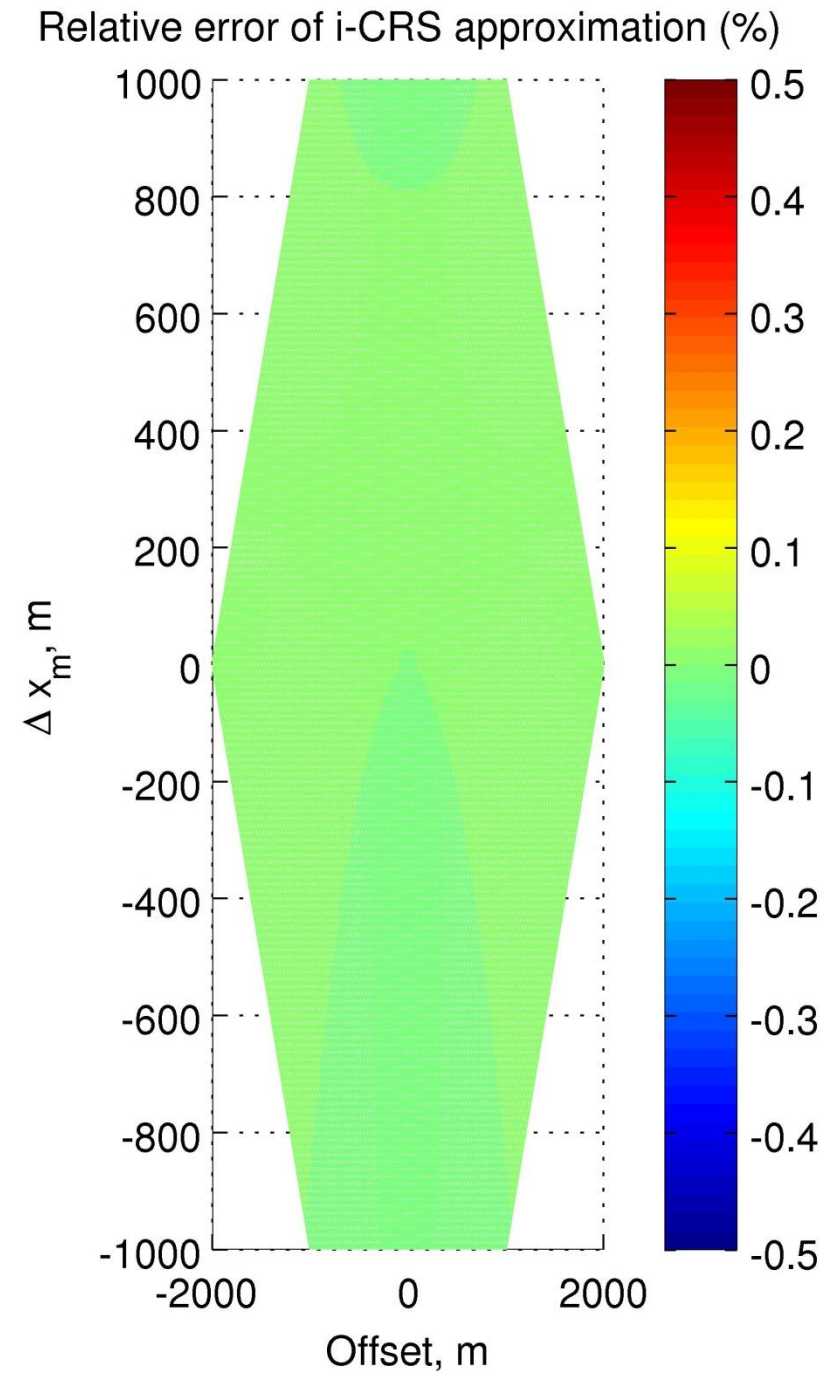
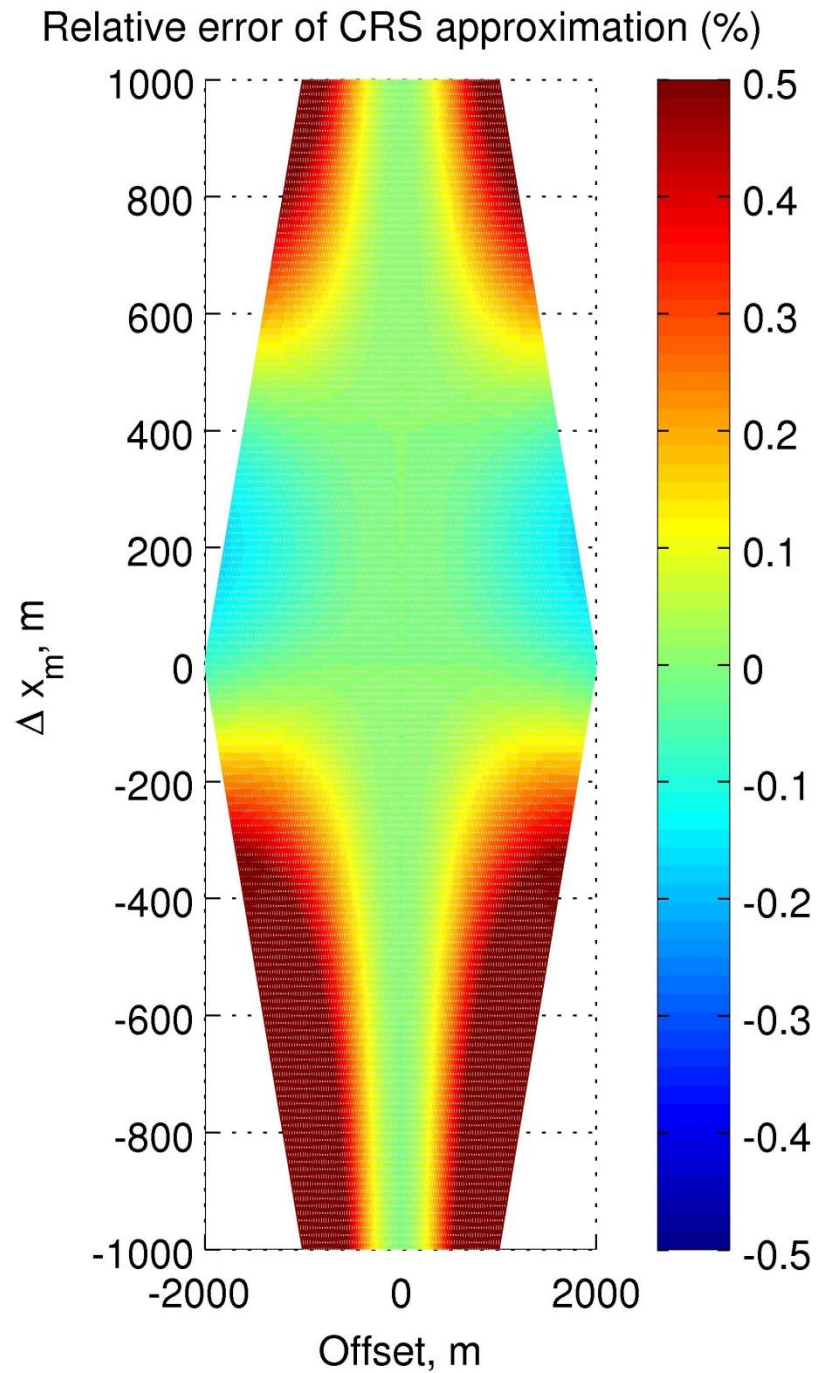
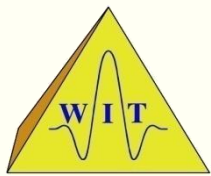


Numerical example: traveltimes

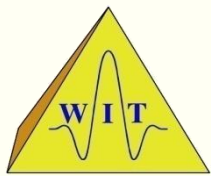


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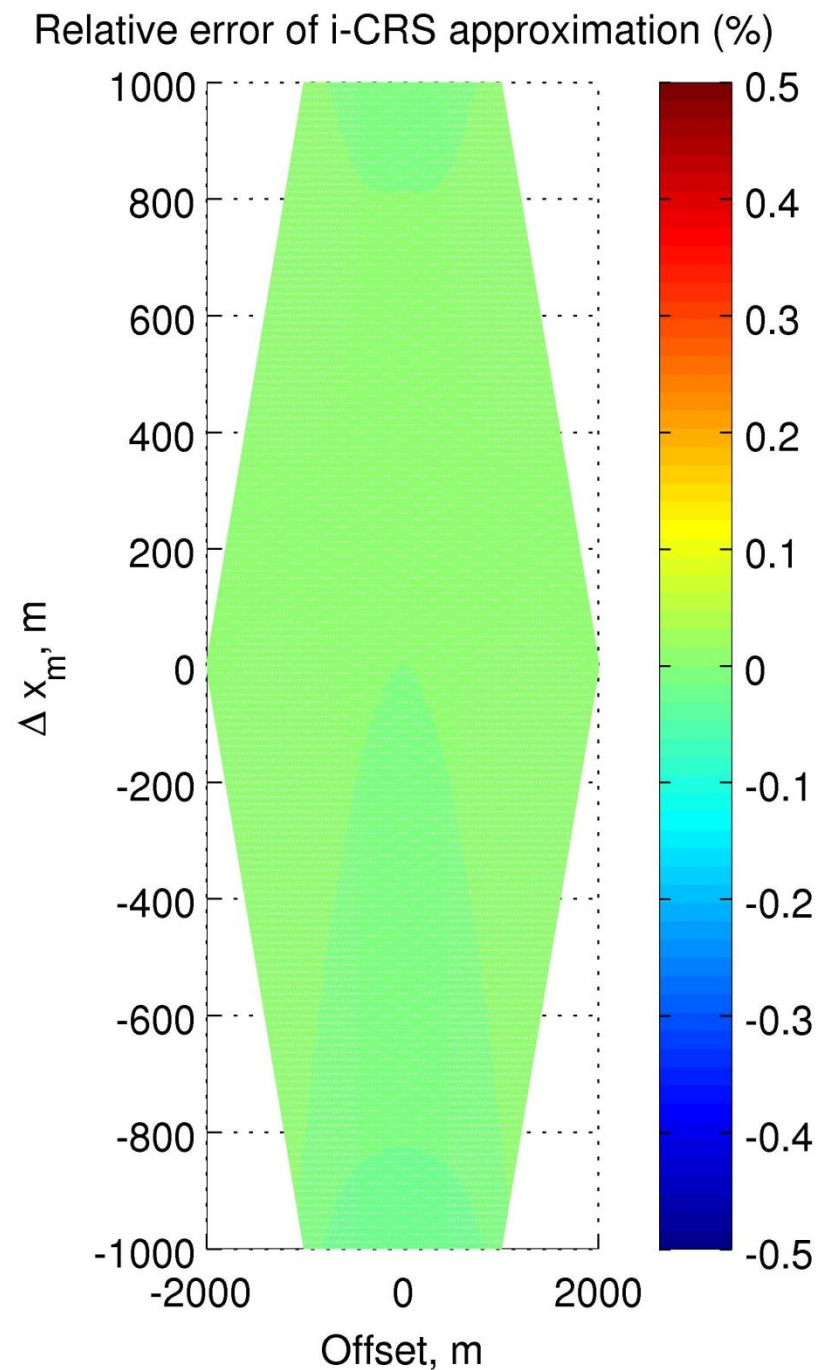
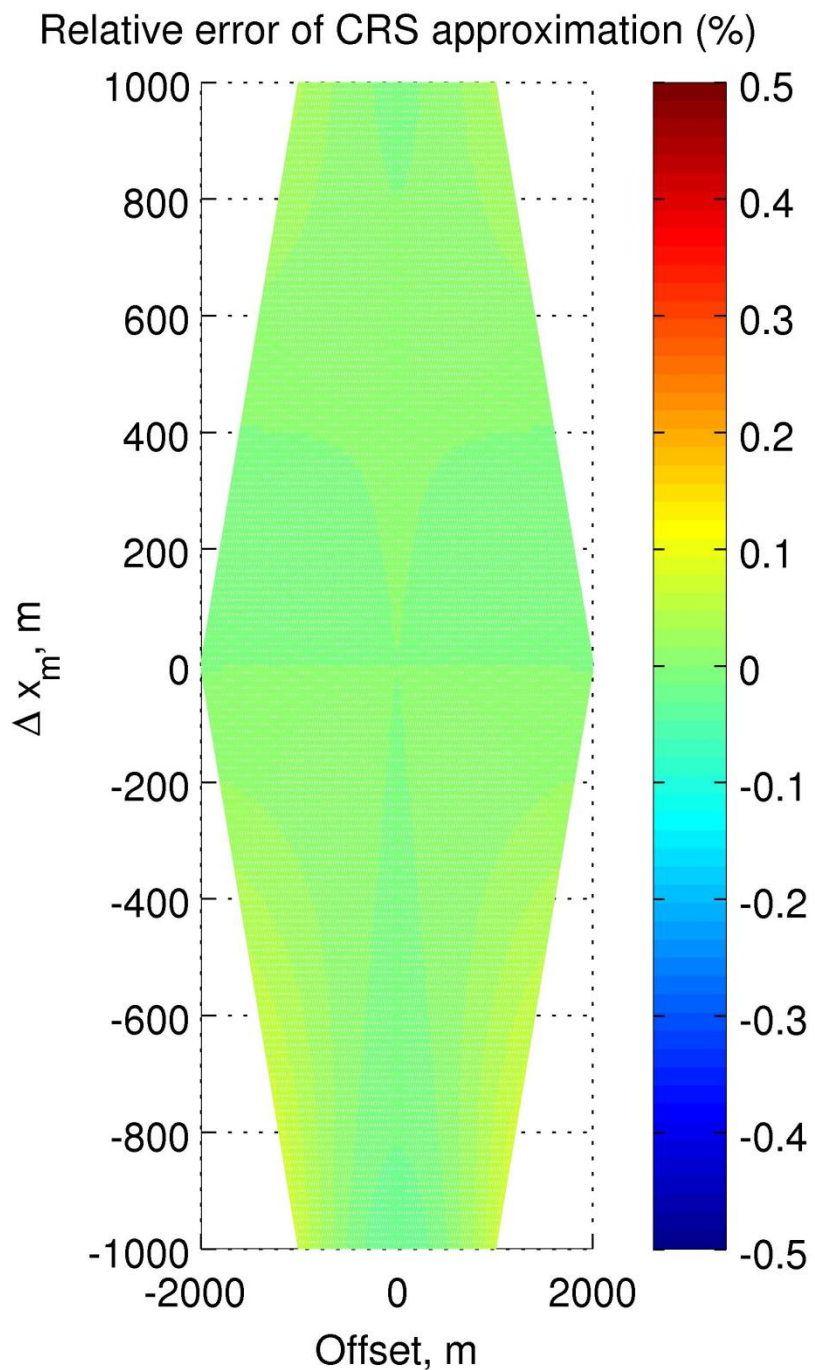
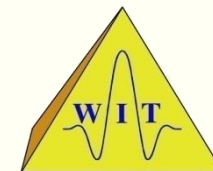


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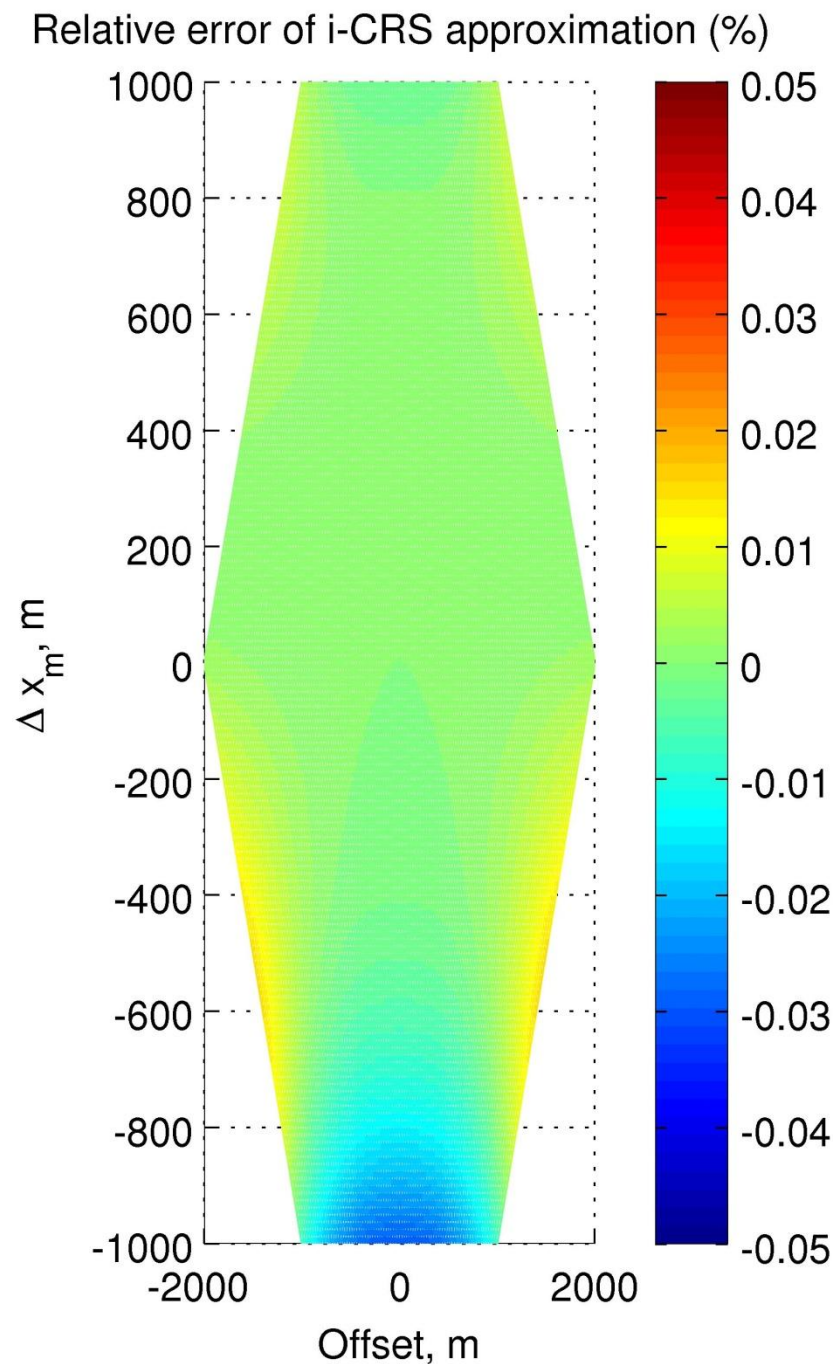
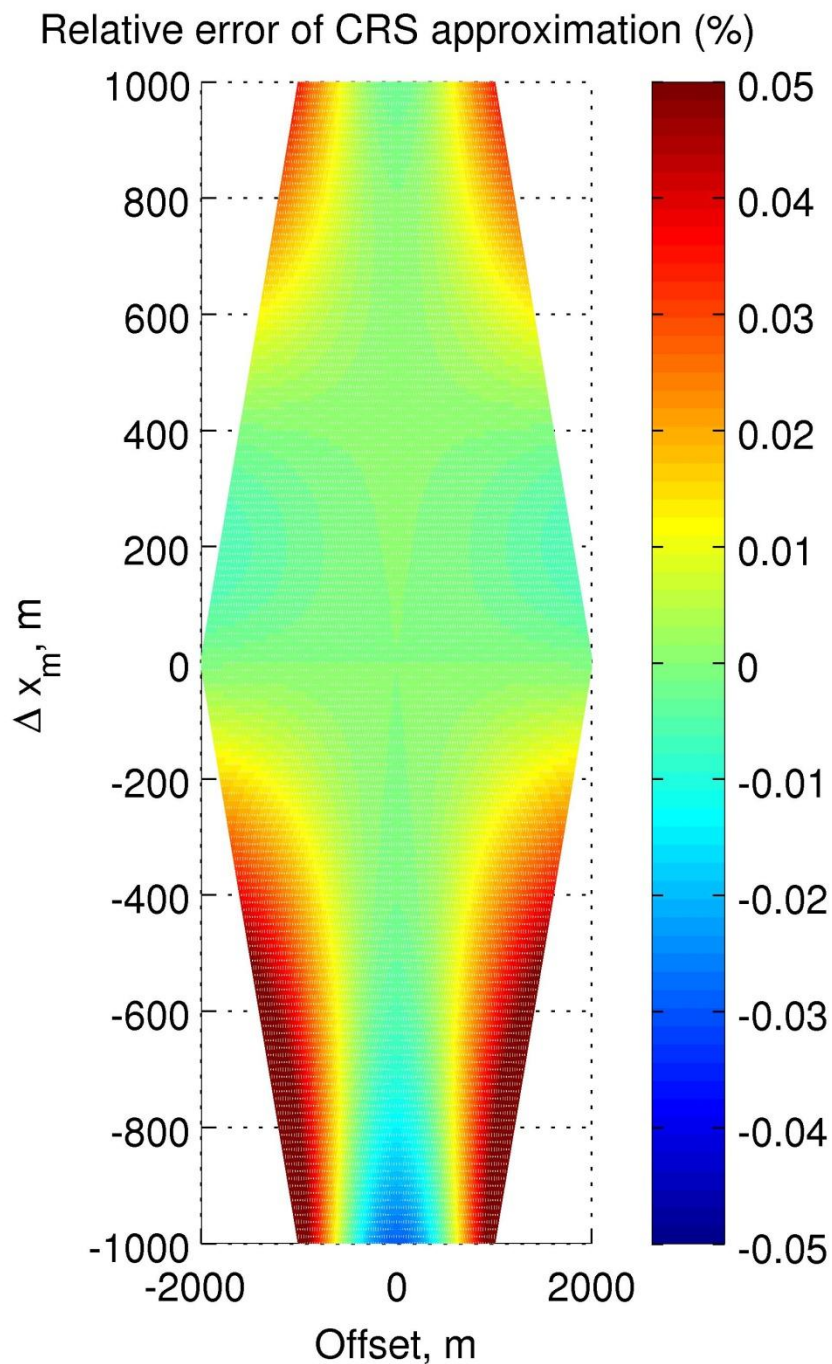
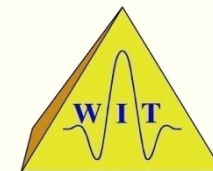


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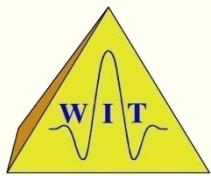
Numerical example: traveltimes



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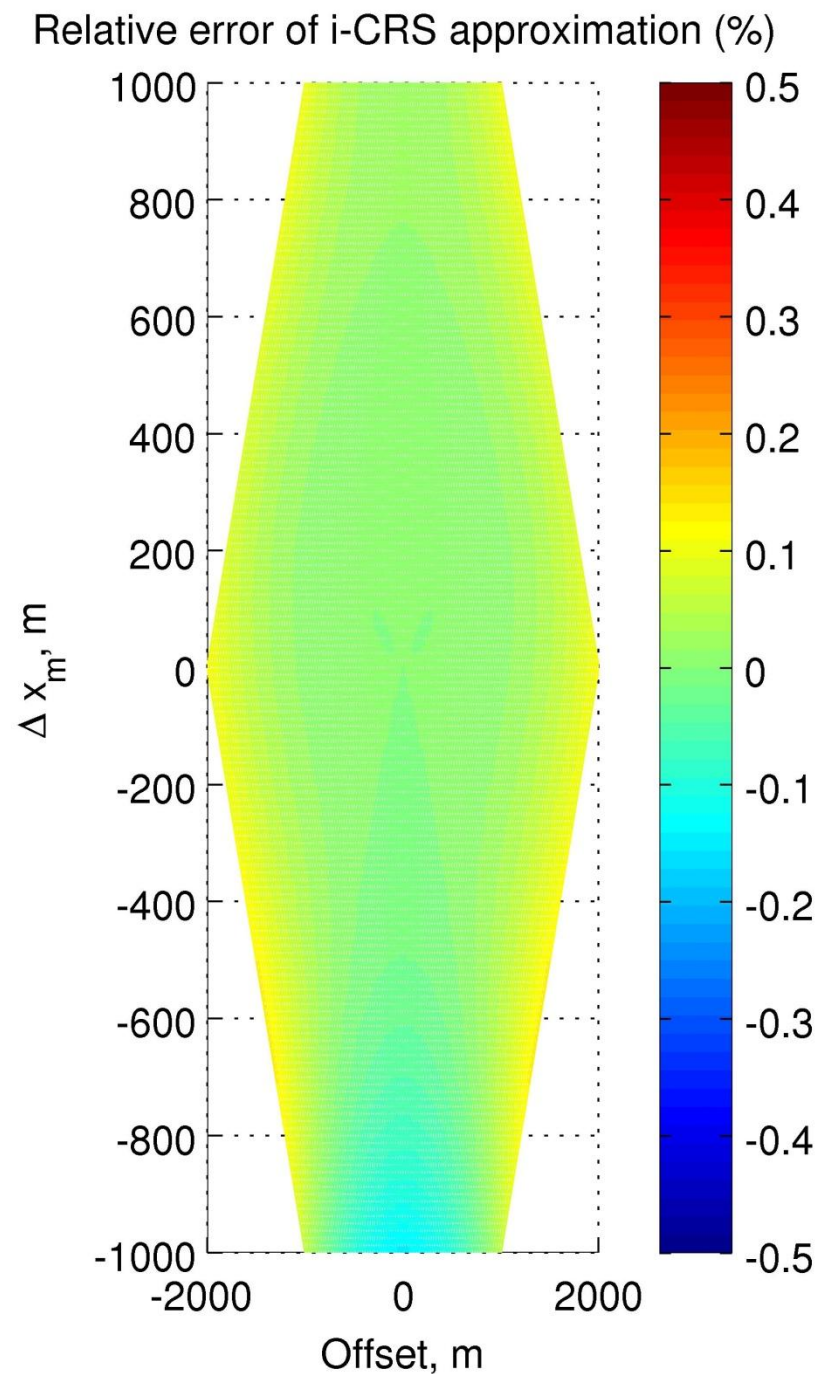
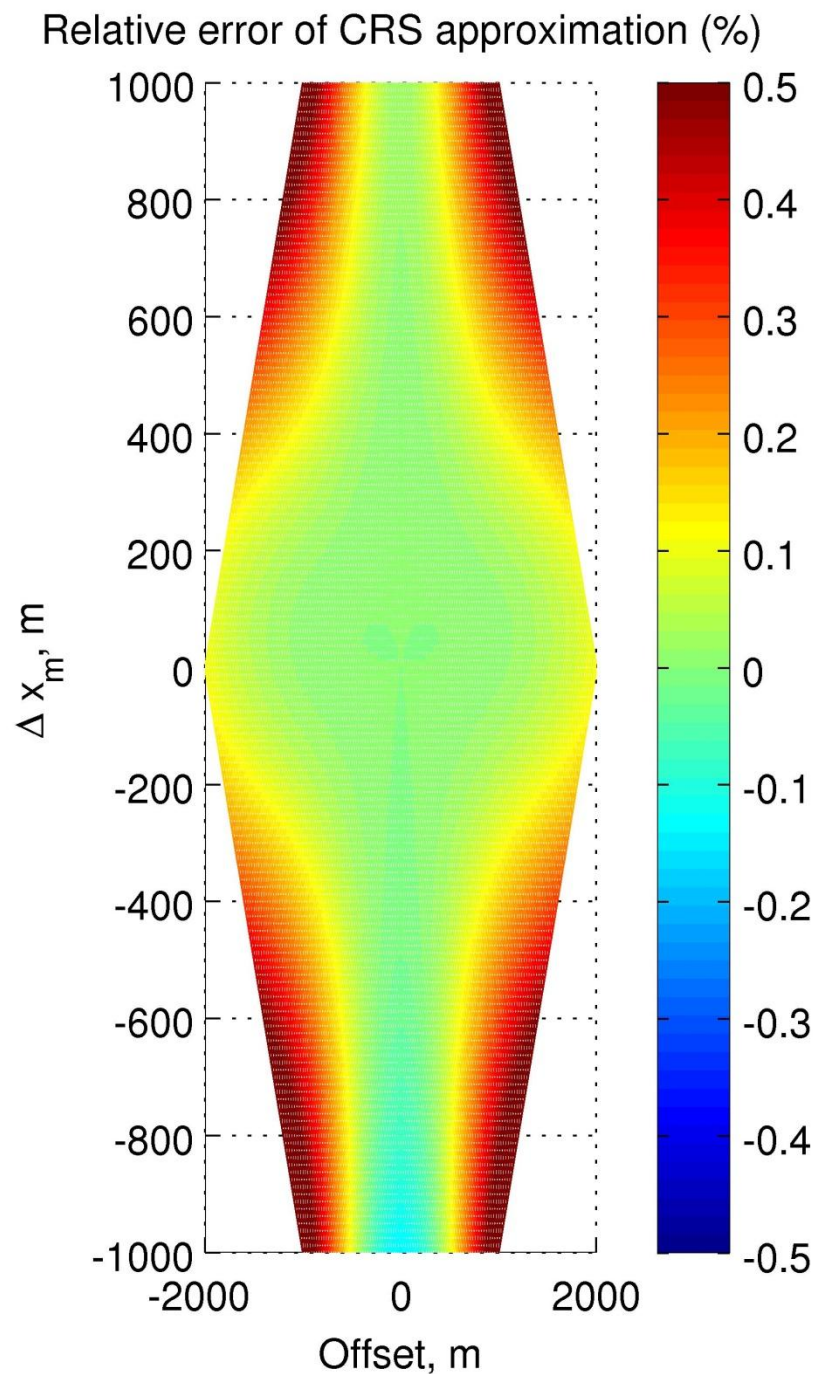
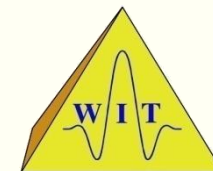


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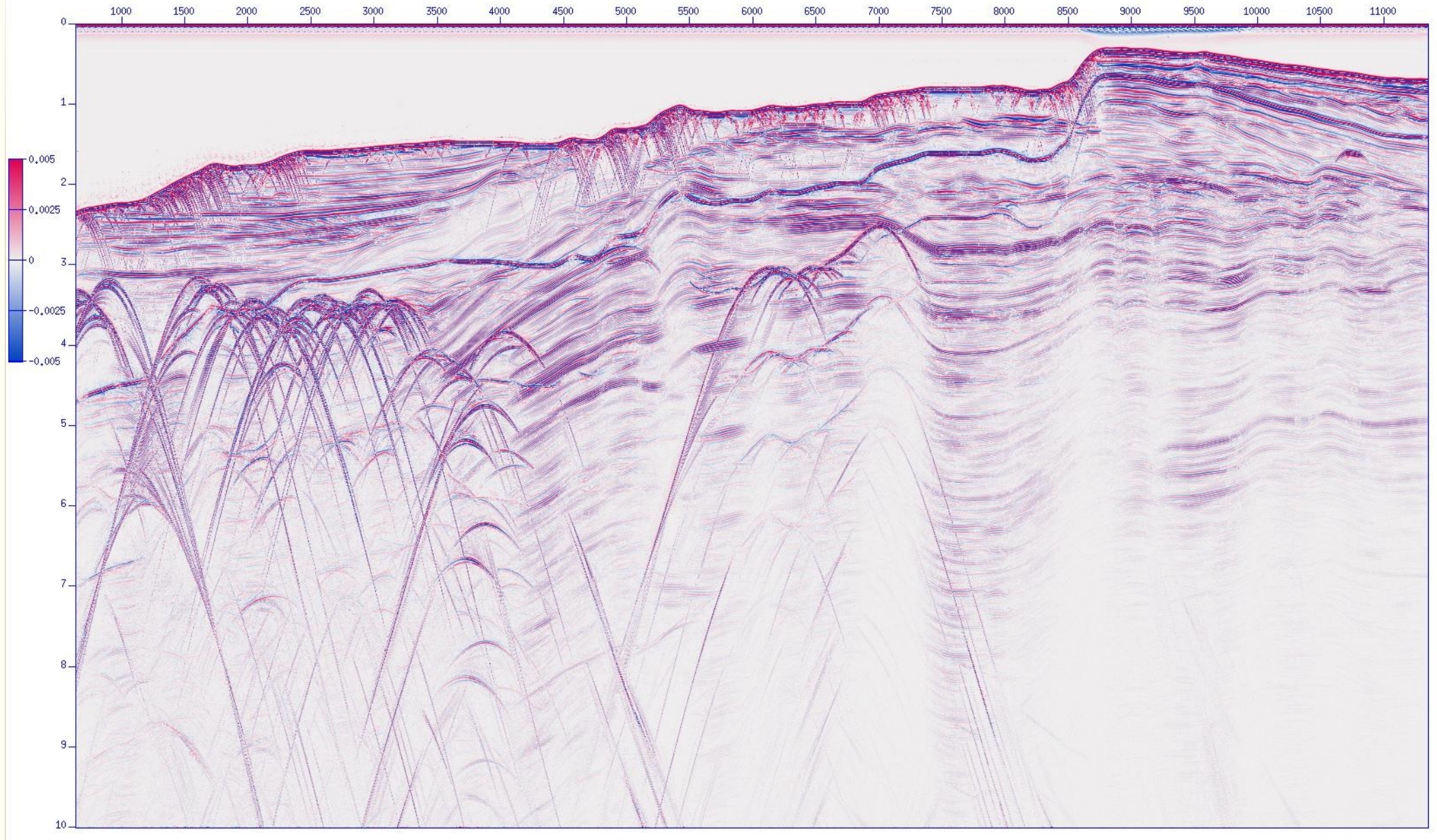
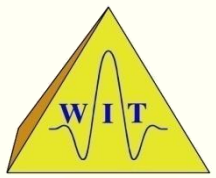


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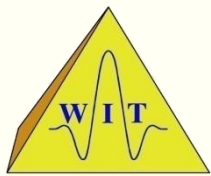
Numerical example: traveltimes



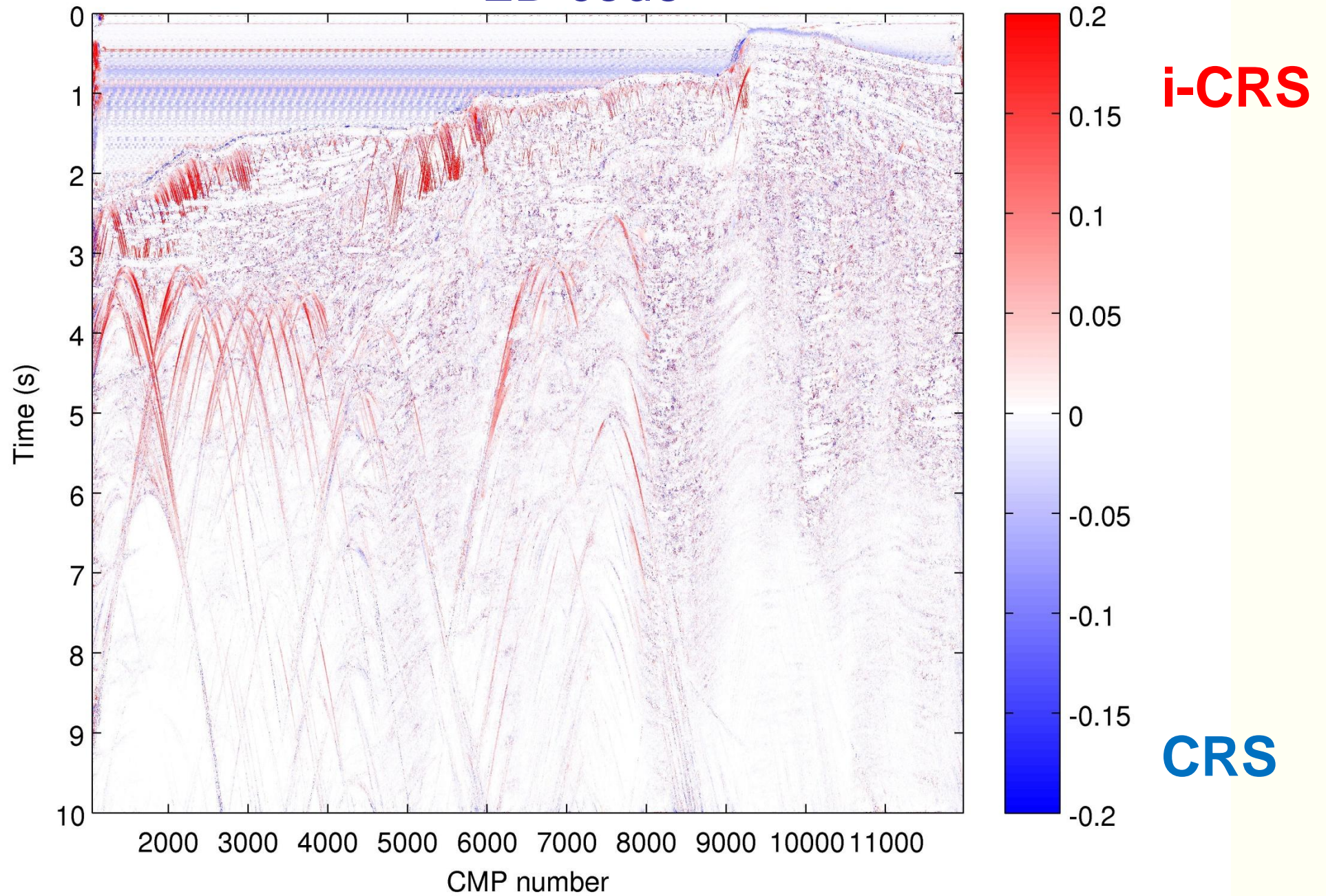
BP test model



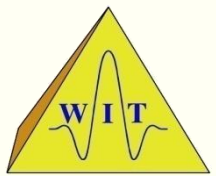
BP test model



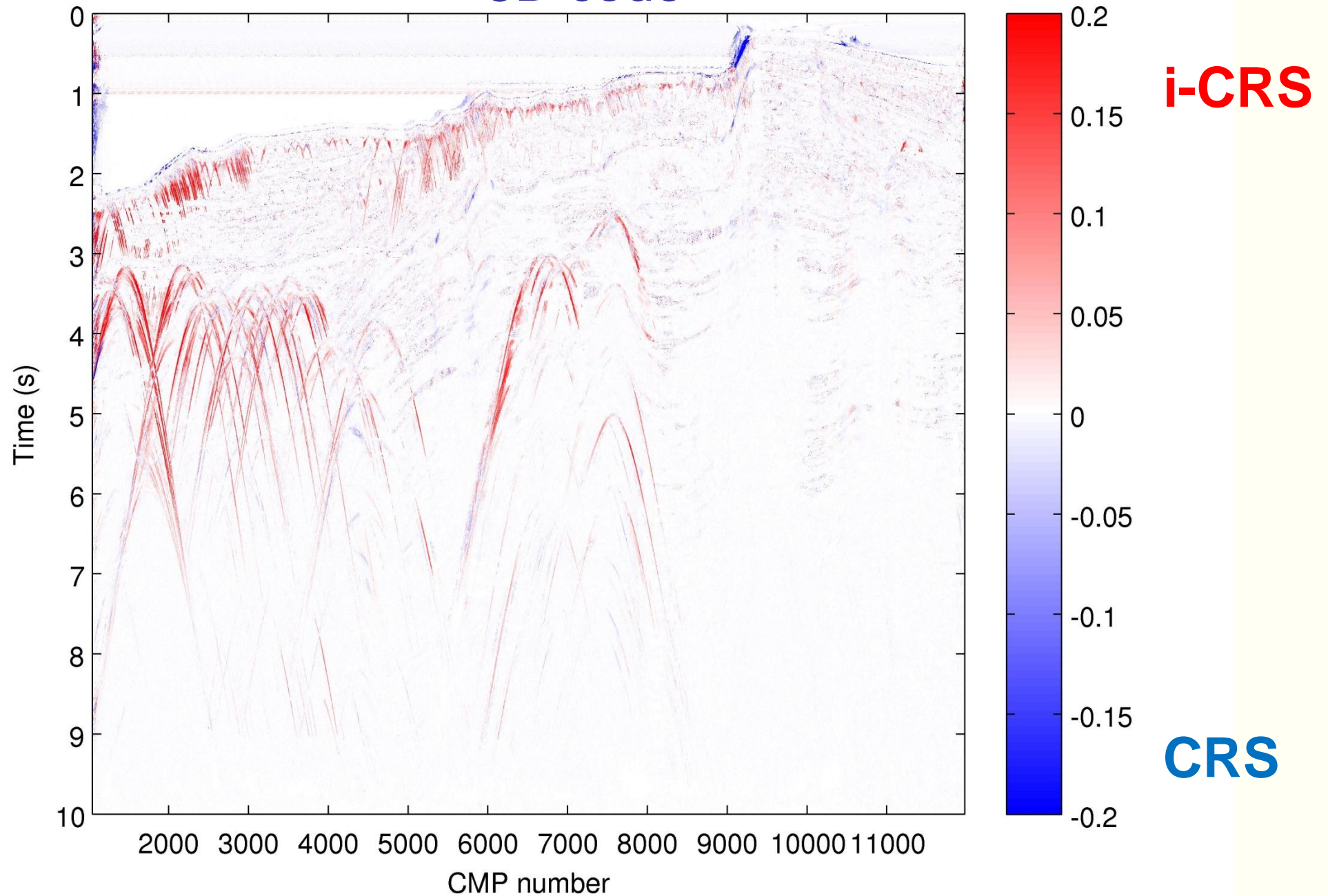
2D code



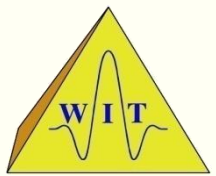
BP test model



3D code



Conclusions



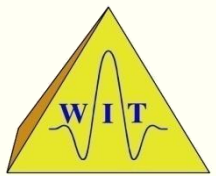
New 3D i-CRS stacking operator:

- intuitive DSR expression
- perfectly fits scattered waves
- easy to implement in the standard CRS workflow
- reasonable price for improved accuracy as compared with standard CRS operator

Effective anisotropic medium and CO geometry:

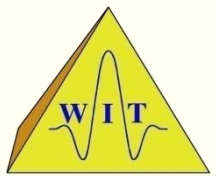
- key point for understanding traveltimes operators in anisotropic medium

Outlook



- Extension to common offset case
- Anisotropy and converted waves
- Improved migration algorithms

Acknowledgments



We would like to warmly thank:

- Sergius Dell (CGG)
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